



Blatt 11 zur Theoretischen Physik IV, WS2023/2024
(Abgabe bis 26.01.2024, 8.30 Uhr)

Exercise 1 *Einstein and Debye model for phonons* [12 Points]

In the following, the contribution of phonons (bosonic quasiparticles of lattice vibrations) to the heat capacity of solids will be investigated. To do this, consider N atoms, each with three isotropic vibration modes. Each vibration mode should represent a harmonic oscillator with frequency ω_i . Analogous to photons, the chemical potential for phonons disappears.

- a) Specify the mean internal energy as a function of the inverse temperature β and the discrete frequencies ω_i . For large N , these frequencies are very close, so that they can also be regarded as continuously distributed. What is the average internal energy?

Hint: Take the density of states $Z(\omega)$ (i.e. in the interval $[\omega, \omega + d\omega]$ there are $Z(\omega) d\omega$ possible frequencies) as given.

In the *Einstein model*, it is now assumed for the sake of simplicity that all oscillators oscillate at the same frequency ω_E .

- b) Use this to explicitly calculate the mean value of the internal energy and the heat capacity. Examine the limiting cases $\beta \rightarrow 0$ and $\beta \rightarrow \infty$ and compare with the Dulong-Petit and Debye's T^3 -law.

The *Debye model* provides an improved description. The possible frequencies are continuously distributed and limited upwards by a specified cut-off frequency ω_D . The density of states is taken to be

$$Z(\omega) = \alpha \omega^2 \Theta(\omega_D - \omega).$$

with $\alpha = \text{const}$ and Θ as a Heaviside function.

- c) How must α be chosen so that the density of states is correctly normalized? Calculate the mean internal energy and the heat capacity again.

Hint: You do not need to solve the integrals for internal energy and heat capacity here.

- d) Now introduce the *Debye temperature* $\Theta_D = \frac{\hbar \omega_D}{k_B}$ and calculate the heat capacity explicitly $T \gg \Theta_D$ or $T \ll \Theta_D$. Compare with the result from b.).

Hint: Develop the numerator and denominator of the integral for $T \gg \Theta_D$ in leading order. Use the following identity for $T \ll \Theta_D$:

$$\int_0^\infty \frac{x^4 e^x}{(e^x - 1)^2} dx = \frac{4\pi^4}{15}.$$

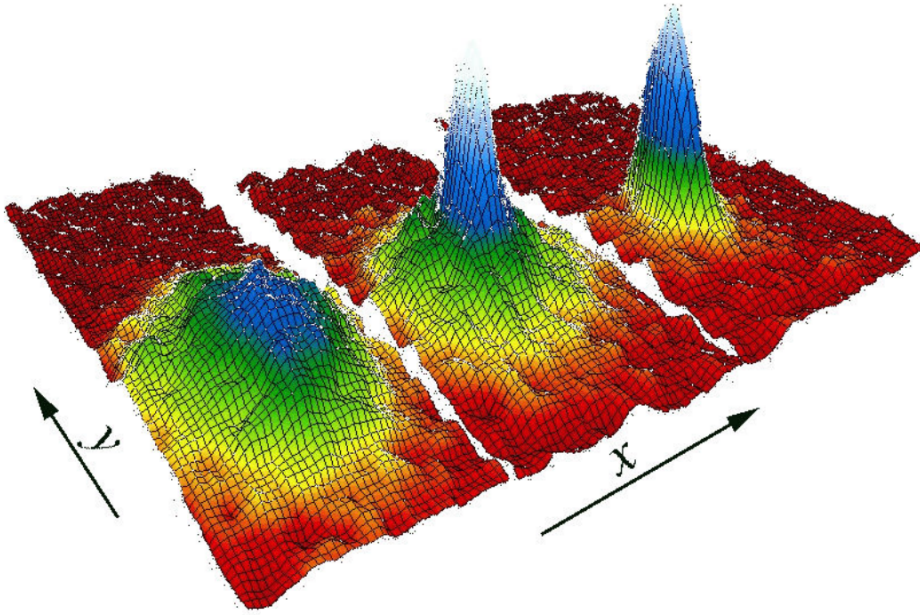


Figure 1: **Bose–Einstein condensation** at 400, 200, and 50 nano-Kelvin. The pictures are spatial distributions 60 ms after the potential is removed; the field of view of each image is $200 \mu\text{m} \times 270 \mu\text{m}$. The left picture is roughly spherically symmetric, and is taken before Bose condensation; the middle has an elliptical Bose condensate superimposed on the spherical thermal background; the right picture is nearly pure condensate. From Anderson, M. H., Ensher, J. R., Matthews, M. R., Wieman, C. E., and Cornell, E. A. (1995). Observation of Bose–Einstein condensation in a dilute atomic vapor. *Science*, 269, 198.

Exercise 2 *Bose condensation: the experiment* [12 Points]

Anderson, Ensher, Matthews, Wieman, and Cornell in 1995 were able to get a dilute gas of rubidium-87 atoms to Bose condense.

- Is rubidium-87 (37 protons and electrons, 50 neutrons) a boson or a fermion?
- At their quoted maximum number density of $2.5 \times 10^{12}/\text{cm}^3$, at what temperature T_c^{predict} do you expect the onset of Bose condensation in free space? They claim that they found Bose condensation starting at a temperature of $T_c^{\text{measured}} = 170 \text{ nK}$. Is that above or below your estimate? (Useful constants: $h = 6.6262 \times 10^{-27} \text{ erg s}$, $m_n \sim m_p = 1.6726 \times 10^{-24} \text{ g}$, $k_B = 1.3807 \times 10^{-16} \text{ erg/K}$.)

The trap had an effective potential energy that was harmonic in the three directions, but anisotropic with cylindrical symmetry. The frequency along the cylindrical axis was $f_0 = 120\text{Hz}$ so $\omega_0 \sim 750\text{Hz}$, and the two other frequencies were smaller by a factor of $\sqrt{8}$: $\omega_1 \sim 265\text{Hz}$. The Bose condensation was observed by abruptly removing the trap potential, and letting the gas atoms spread out; the spreading cloud was imaged 60 ms later by shining a laser on them and using a CCD to image the shadow.

For your convenience, the ground state of a particle of mass m in a one-dimensional harmonic oscillator with frequency ω is $\psi_0(x) = (\frac{m\omega}{\pi\hbar})^{1/4} e^{-m\omega x^2/2\hbar}$, and the momentum-space wavefunction is $\tilde{\psi}_0(p) = (\pi m\hbar\omega)^{-1/4} e^{-p^2/2m\hbar\omega}$. In this 3D problem the solution is a product of the corresponding Gaussians along the three axes.

- Will the momentum distribution be broader along the high-frequency axis (ω_0) or one of the low-frequency axes (ω_1)?
Assume that you may ignore the small width in the initial position distribution, and that the positions in Fig. 1 thus reflect the velocity distribution times the time elapsed. Which axis, x or y in Fig. 1, corresponds to the high-frequency cylinder axis? What anisotropy would one expect in the momentum distribution at high temperatures (classical statistical mechanics)? Does this high-temperature anisotropy agree with the experimental measurement in 1?

Their Bose condensation is not in free space; the atoms are in a harmonic oscillator potential. In the calculation in free space, we approximated the quantum states as a continuum density of states $g(E)$. That is only sensible if $k_B T$ is large compared to the level spacing near the ground state.

- d) Compare $\hbar\omega$ to $k_B T$ at the Bose condensation point T_c measured in their experiment. ($\hbar = 1.05459 \times 10^{-27}$ erg s; $k_B = 1.3807 \times 10^{-16}$ erg/K.)

For bosons in a one-dimensional harmonic oscillator of frequency ω_0 , it is clear that $g(E) = 1/(\hbar\omega_0)$; the number of states in a small range ΔE is the number of $\hbar\omega_0$ it contains.

- e) Calculate the density of single-particle eigenstates

$$g(E) = \int_0^\infty d\epsilon_1 d\epsilon_2 d\epsilon_3 g_1(\epsilon_1) g_2(\epsilon_2) g_3(\epsilon_3) \times \delta(E - (\epsilon_1 + \epsilon_2 + \epsilon_3))$$

for a three-dimensional harmonic oscillator, with one frequency ω_0 and two of frequency ω_1 .

Hint: What is the shape of the energy shell in ϵ space?

Their experiment has $N = 2 \times 10^4$ atoms in the trap as it condenses.

- f) By working in analogy with the calculation in free space and your density of states from part (e), find the maximum number of atoms that can occupy the three-dimensional harmonic oscillator potential (without Bose condensation) at temperature T .

Hint: $\int_0^\infty z^2/(e^z - 1) dz = 2\zeta(3) = 2.40411$.

According to your calculation, at what temperature T_c^{H0} should the real experimental trap have Bose condensed?

Exercise 3 One-dimensional XY-model [12 Points]

Consider a 1D spin model, which is defined by the variables ϕ_i , each of which can take the values $-\pi$ and π and interact with their neighbors ϕ_{i+1} and ϕ_{i-1} such that the sum of states is given by

$$Z = \int_{-\pi}^{\pi} \prod_{l=1}^N \frac{d\phi_l}{2\pi} e^{K \sum_{k=1}^{N-1} \cos(\phi_k - \phi_{k+1}) + K \cos(\phi_N - \phi_1)}.$$

- a) Find the transfer matrix.
b) Calculate the sum of states in the limit $N \rightarrow \infty$.

Hint: The eigenvalue equation is an integral equation and the normalized eigenfunctions of the transfer matrix are $\frac{e^{in\phi}}{\sqrt{2\pi}}$ for $n = 0, \pm 1, \pm 2, \dots$. For real a and integer n , you can use the modified Bessel function $I_n(a) = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{a \cos \phi + in\phi} d\phi$.

- c) Determine the correlation length ξ in the limiting case $N \rightarrow \infty$.

Exercise 4 Computational task: 2D Ising model of ferromagnetism [4 Points]

Simulate the 2D Ising model of ferromagnetism, and calculate the magnetization of the system at different temperatures. Interpret the result. As a guide, there is a Python code in the exercise folder (blatt11_2D-Ising.py), you can complete it.